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Spherical Waves

The utility of thinking of \vec{k} as a "ray" becomes even more obvious when we get away from plane waves and start thinking of waves with *curved* wavefronts. The simplest such wave is the type that is emitted when a pebble is tossed into a still pool - an example of the "point source" that radiates waves isotropically in all directions. The wavefronts are then *circles* in two dimensions (the surface of the pool) or *spheres* in three dimensions (as for sound waves) separated by one wavelength λ and heading outward from the source at the propagation velocity c . In this case the "rays" \vec{k} point along the radius vector \hat{r} from the source at any position and we can once again write down a rather simple formula for the "wave function" (displacement A as a function of position) that depends only on the time t and the *scalar* distance r from the source.

A plausible first guess would be just $A(x, t) = A_0 e^{i(kr - \omega t)}$, but this cannot be right! Why not? Because it violates energy conservation. The energy density stored in a wave is proportional to the square of its amplitude; in the trial solution above, the amplitude of the outgoing spherical wavefront is constant as a function of r , but the *area* of that wavefront increases as r^2 . Thus the energy in the wavefront increases as r^2 ? I think not. We can get rid of this effect by just dividing the amplitude by r (which divides the energy density by r^2). Thus a trial solution is

$$A(x, t) = A_0 \frac{e^{i(kr - \omega t)}}{r} \quad (14.38)$$

which is, as usual, correct.^{14.11} The factor of $1/r$ accounts for the conservation of energy in the outgoing wave: since the spherical "wave front" distributes the wave's energy over a surface area $4\pi r^2$ and the flux of energy per unit area through a spherical surface of radius r is proportional to the *square* of the wave amplitude at that radius, the integral of $|f|^2$ over the entire sphere (*i.e.* the total outgoing *power*) is independent of r , as it must be.

We won't use this equation for anything right now, but it is interesting to know that it does accurately describe an outgoing^{14.12} spherical wave.

The perceptive reader will have noticed by now that Eq. (38) is not a solution to the WAVE EQUATION as represented in one dimension by Eq. (10). That is hardly surprising, since the spherical wave solution is an intrinsically 3-dimensional beast; what happened to y and z ? The correct *vector* form of the WAVE EQUATION is

$$\nabla^2 A - \frac{1}{c^2} \frac{\partial^2 A}{\partial t^2} = 0 \quad (14.39)$$

where the LAPLACIAN operator ∇^2 can be expressed in Cartesian^{14.13} coordinates (x, y, z) as^{14.14}

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \quad (14.40)$$

With a little patient effort you can show that Eq. (38) does indeed satisfy Eq. (39), if you remember that $r = \sqrt{x^2 + y^2 + z^2}$. Or you can just take my word for it . . .

Footnotes

. . . correct. [14.11](#)

I should probably show you a few wrong guesses first, just to avoid giving the false impression that we always guess right the first time in Physics; but it would use up a lot of space for little purpose; and besides, "knowing the answer" is always the most powerful problem-solving technique!

. . . outgoing [14.12](#)

One can also have "incoming" spherical waves, for which Eq. (38) becomes

$$A(x, t) = A_0 \frac{e^{i(kr + \omega t)}}{r}.$$

. . . Cartesian [14.13](#)

The LAPLACIAN operator can also be represented in other coordinate systems such as spherical (r, θ, ϕ) or cylindrical (ρ, θ, z) coordinates, but I won't get carried away here.

. . . as [14.14](#)

The LAPLACIAN operator can also be thought of as the inner (scalar or "dot") product of the GRADIENT operator $\vec{\nabla}$ with itself: $\nabla^2 = \vec{\nabla} \cdot \vec{\nabla}$, where

$$\vec{\nabla} = \hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z}$$

in Cartesian coordinates. This VECTOR CALCULUS stuff is really elegant - you should check it out sometime - but it is usually regarded to be beyond the scope of an introductory presentation like this.

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Jess H. Brewer 2002-03-26